Vibrational spectra study of $Mo(V)$ phosphates as examples of different geometries of dimolybdenyl species[†]

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Received 13th September 2000, Accepted 22nd March 2001 First published as an Advance Article on the web 26th April 2001

Three different phosphate structures, α -K₂Mo₂P₂O₁₁, β -K₂Mo₂P₂O₁₁ and CdMoPO₆, in which molybdenum(v) octahedra are assembled in different ways, have been studied by vibrational spectroscopy techniques. A correlation is given between crystal structure and infrared and Raman band behaviour which allows identification of vibrational modes sensitive to a particular configuration of molybdenum species. This represents a powerful tool for structure characterisation when amorphous or single layer supported phases are under study.

Introduction

Powders based on molybdenum oxides are the object of great interest because of their excellent behaviour as catalysts for selective oxidation.¹ Mixing of transition metal oxides with other oxides is frequently helpful to improve activity/selectivity as well as to obtain optimal acid–basic behaviour. Phosphoric anhydride is frequently used as a promoter component for selective oxidation catalysts, so that Mo–phosphates, either supported or as bulk materials, are reported to be interesting catalytic materials.^{2–5} Mo(v) phosphates can be considered as analogues of $V(w)$ phosphates which represent the active phases for n-butane selective oxidation catalysts.^{6,7} These catalysts are also expected to display quite strong acidic properties,^{5,8,9} that can be moderated by doping with alkali metal ions such as potassium. For these reasons, the structure and spectroscopic properties of well characterised K–Mo–P mixed oxides are of interest as reference materials for the characterisation of supported heterogeneous catalysts with similar compositions.

Recently, several new phases of phosphates of molybdenum(v_I) and molybdenum(v) have been prepared and characterised.^{10–17} In particular, two different allotropic forms of potassium–molybdenum(v) phosphates, i.e. α -K₂Mo₂P₂O₁₁ and β -K₂Mo₂P₂O₁₁, have been synthesized and characterized.^{11,15,16} They present different structural units of $MoO₆$ octahedra and \overline{PO}_4 tetrahedra: corner-shared pairs of MoO_6 octahedra occur in the monophosphate α -K₂Mo₂P₂O₁₁, while edge-shared pairs of $MoO₆$ octahedra occur in the diphosphate β -K₂Mo₂P₂O₁₁. For further comparison, we have also investigated another $Mo(v)$ monophosphate, $CdMoPO₆,¹⁷$ which also contains edge-shared pairs of $MoO₆$ octahedra. The features of this phase will be compared with those of the $CdMo^{VI}O₄ phase, previously investigated.¹⁴ We report here an$ IR and Raman study of these three pentavalent molybdenum phosphates, the study being presented taking into account the increasing complexity of molybdo–phosphate units: first the phase constituted by octahedra sharing one edge is investigated, then that containing octahedra sharing one corner and connected by two phosphates, and finally $CdMoPO₆$ where two bioctahedral units are linked by two phosphates.

Experimental

 α -K₂Mo₂P₂O₁₁, β -K₂Mo₂P₂O₁₁ and CdMoPO₆ powders were obtained via solid state reaction of precursors in an evacuated sealed silica tube at 973, 923 and 1173 K, respectively. Details relating to synthesis and phase recognition are reported in the literature.^{11,15–17}

The IR and FIR spectra were recorded on a Nicolet Magna 750 Fourier transform instrument. The FT-Raman spectra were recorded using a Brucker FTS100 (Nd-YAG laser). The skeletal spectra in the region above 400 cm^{-1} were recorded with KBr pressed disks and with a KBr beam splitter, while those in the far infrared region $(400-50 \text{ cm}^{-1})$ were recorded using powder deposited on polyethylene disks, and with a ''solid substrate'' beam splitter.

Results and discussion

1. The vibrational spectra of β -K₂Mo₂P₂O₁₁

For the phase β -K₂Mo₂P₂O₁₁ we could record high quality FT-IR/FT-FIR and FT-Raman spectra, which are shown in Fig. 1a and 2a, respectively.

As determined previously by Guesdon et al. in their XRD study on single crystals and powders,^{11,16} β -K₂Mo₂P₂O₁₁ crystallises in the *Pbcn* space group (D_{2h}^{14} , no. 60) with $Z=4$. It is a diphosphate, containing a complex $Mo(v)$ oxocation, with formula $[Mo₂O₄]²⁺$, so that this compound is better described by the formula $K_2[M_02O_4][P_2O_7]$. Consequently, it can be considered as an allotropic form with respect to the corresponding high-temperature form α -K₂Mo₂P₂O₁₁,^{11,15} which is a monophosphate and can be expressed by the formula $K_2Mo_2O_3[PO_4]_2$.

The coordination around molybdenum in β -K₂Mo₂P₂O₁₁ is shown schematically in Fig. 3a. It presents $Mo₂O₁₀$ bioctahedral units, formed by two edge-sharing octahedra. The Mo–O–Mo angle is 83.10° . Each of the octahedra around Mo has three

[{]Electronic supplementary information (ESI) available: determination of the irreducible representation. and separation into internal and external modes for β -K₂Mo₂P₂O₁₁. See http://www.rsc.org/suppdata/ jm/b0/b007416k/

Fig. 1 FT-IR/FT-FIR spectra of β -K₂Mo₂P₂O₁₁ (a), α -K₂Mo₂P₂O₁₁ (b) and CdMoPO₆ (c). The spectrum in the FIR region was recorded by deposition of the pure powder on a polyethylene substrate.

coordinated diphosphate groups, two bridging oxygen atoms with the nearer Mo octahedron and one free corner or, more precisely, one short Mo=O "molybdenyl" bond (1.697 Å) . This abnormally short bond is a characteristic feature of $Mo(V)O₆$ octahedra. The corresponding terminal oxygen atom interacts in this structure with two K atoms. The two molybdenyl bonds of the $Mo₂O₁₀$ bioctahedral unit can be considered as essentially parallel to each other and to the **axis, although they are slightly** divergent. The "cyclic" $Mo₂O₂$ units, nearly parallel to the *ac* plane, are not exactly square planar, the oxygen atoms lying at a lower level with respect to the molybdenum atoms. The diphosphate group has a bent P–O–P bridge (P–O–P 138.2 $^{\circ}$) and a "random" conformation of the $PO₃$ units.

On the basis of the site symmetries of the positions occupied in this structure^{11,16} and according to the correlation method,¹⁸ factor group analysis allowed us to obtain the following irreducible representation:

$$
\begin{aligned} \varGamma_{opt} = & 28A_g + 29B_{1g} + 28B_{2g} + 29B_{3g} + 28A_u \\ & + 29B_{1u} + 28B_{2u} + 29B_{3u} \end{aligned}
$$

According to the structure described above, we can separate the total modes of the structure into those associated to the diphosphate groups, to the $Mo₂O₄$ complex ion and to potassium, and we can also separate the vibrational and the acoustic modes. This is shown in Table 1, where the internal modes forecasted for the diphosphate ions and for the $Mo₂O₄$ complex cations are also identified. In Table 2 the observed maxima in the IR and Raman spectra of Figs. 1 and 2 are listed and assignment for the higher frequency features attempted.

In the fundamental IR spectrum (\bar{v} < 1300 cm⁻¹) we note the presence of 44 features among maxima and shoulders, while in the Raman spectra 13 maxima and several shoulders can be found. Thus many less features with respect to the number of expected IR-active and Raman active modes (77 and 102, respectively, see Table 1) are actually found.

In the region $1300-1000 \text{ cm}^{-1}$ of the IR spectrum we distinguish eight features, between maxima and pronounced shoulders. This region is typical for the asymmetric and symmetric stretches of terminal PO bonds, in this case $v_{as}(PO_3)$ and $v_s(PO_3)$. According to the structure we forecast 9 IR-active modes, which means that only one of these eludes detection, probably because it is superimposed on another. In the Raman spectrum, in the same region we only find three well-resolved maxima, in contrast with the 12 Raman-active expected modes. However, we can remark that among these modes, three have A_{α} symmetry, *i.e.* are totally symmetric. The A_{α} symmetry modes are expected to be the most intense in the Raman spectrum. Consequently, we can assign the three Raman peaks

Fig. 2 Raman spectra of β -K₂Mo₂P₂O₁₁ (a) and CdMoPO₆ (c) in the range of the vibrational modes.

Fig. 3 The oxygen coordination around molybdenum and phosphorus seen from different directions for: β -K₂Mo₂P₂O₁₁ (a), α -K₂Mo₂P₂O₁₁ (b) and $CdMoPO₆$ (c).

to the three A_g symmetry modes, one of which is a symmetric stretch and two are asymmetric stretches.

In the region $1000-800 \text{ cm}^{-1}$ of the IR spectrum we find a complex band whose maximum is definitely split but contains at least one shoulder at the higher frequency side. In the Raman spectrum a strong sharp maximum is found, with at least two additional weak components, one for each side. In this region, the asymmetric stretching modes of the P–O–P bridge of diphosphate groups and the stretches of terminal $Mo=O$ bonds can be expected. According to Table 1 we expect here 5 IR active modes and 6 Raman active modes, one only of which belongs to the A_g symmetry. The IR and Raman spectra of different metal diphosphates^{19,23} show that the $v_{\text{as}}(POP)$ modes on pyro- and poly-phosphates are responsible for strong and rather broad IR bands centred in the range $970-940$ cm⁻¹, while the corresponding Raman-active modes are weak and become even inactive when the P–O–P bridges are linear. By contrast, the $Mo=O$ stretching modes are generally strong and sharp both in IR and in Raman. These considerations allow us to confidently assign the strong Raman peak at 948 cm^{-1} to the totally symmetric stretching mode of molybdenyl groups. This position agrees with that of the $Mo=O$ stretch in the Raman spectrum of α -KMoO(P₂O₇), *i.e.* 947 cm⁻¹.²⁴ In the IR spectrum we can suppose that the overall absorption in this region (centred near 950 cm⁻¹) is due to the IR active v_{as} (POP) modes, while the

Table 1 Assignments of the modes associated to P₂O₇ and Mo₂O₄ moieties (4 per unit cell) and to K atoms (8 per unit cell) in β -K₂Mo₂P_{2O11}

Symmetry of modes \rightarrow			A_g	$\mathbf{B}_{1\mathrm{g}}$	$\mathbf{B}_{2\mathbf{g}}$	$\mathbf{B}_{3\mathrm{g}}$	\mathbf{A}_u	$\mathbf{B}_{1\mathrm{u}}$	$\mathbf{B}_{2\mathbf{u}}$	$\mathbf{B}_{3\mathrm{u}}$	
notation \downarrow	\boldsymbol{N}	Act.:	\mathbb{R}	\mathbb{R}	$\mathbf R$	$\mathbf R$	$=$	IR	IR	$\ensuremath{\mathsf{IR}}\xspace$	Total
$v_{\rm as}(\rm PO_3)$			\overline{c}	2	$\overline{2}$	$\overline{2}$	2	2	2	2	16
$v_s(PO_3)$											8
$\delta_{as}(\overline{PO}_3)$				$\overline{2}$				2		2	16
δ _s (PO ₃)											8
Rock PO ₃								2	2	2	16
Twist $PO3$											
$v_{as}(P-O-P)$											
$v_s(P-O-P)$								0			
δ (P-O-P)				Ω		Ω		0			
Internal	21		11	10	11	10	11	10	11	10	84
Translations	3			\overline{c}		\overline{c}		\overline{c}		2	12
Rotations				$\mathfrak{2}$		$\mathfrak{2}$		$\mathfrak{2}$		$\mathfrak{2}$	12
Total	27		13	14	13	14	13	14	13	14	108
$v_{\rm as}(\rm Mo\text{=}O)$			0		θ		θ		Ω		
$v_s(Mo=O)$				0		0		0			
$\delta(Mo=O)$				2	2	2		2	2		16
$v_s/\delta(Mo-O-Mo)$											8
$v_{as}(Mo-O-Mo)$											
v_s (O-Mo-O)											
Rock Mo-O-Mo				Ω		0		0			
Internal	12							5			48
Translations	3			\mathfrak{D}		2		2		2	12
Rotations	3			$\overline{2}$		2		\overline{c}		2	12
Total	18			9		9		9		9	72
K atoms	3			3	3	3	3	3	3	3	24
Total			25	26	25	26	25	27	26	27	204
Vibrational			25	26	25	26	25	26	25	26	201
Acoustic			θ	θ	θ	θ	$\mathbf{0}$	1	1	1	3
N = number of modes per diphosphate ion or $Mo2O4$ molecular unit; Act. = activity.											

two sharp maxima at 936 cm⁻¹ and 928 cm⁻¹ are associated to the IR active $v_s(Mo=O)$ (B_{2u}) and $v_{as}(Mo=O)$ (B_{1u}, B_{3u}) modes.

According to these data, we can assign the weak shoulders in the Raman patterns near 960 cm^{-1} and at 930 cm^{-1} to the Raman active components of the $v_{as}(\text{POP})$ and $v_{as}(\text{Mo}=O)$ modes. The relative intensities of the observed maxima agree with these assignments.

Strong support for these assignments for the absorptions in the $1000-800$ cm⁻¹ region can be obtained by the analysis of the IR spectra in the corresponding first harmonic region, $2000-1600 \text{ cm}^{-1}$ (Fig. 4). Here we found a quite broad band with some components, whose maximum is found at 1896 cm⁻¹. Consisting of first overtones and all self-combinations of either the IR-active or Raman-active modes IRforbidden in the D_{2h} factor group, this band can be assigned to the mixed combinations of the IR- and Raman-active components of the asymmetric stretch of the P–O–P bridge $(\sim 960 + \sim 950 = \sim 1910)$. At lower frequency we found a sharp band with a split maximum $(1872, 1862 \text{ cm}^{-1})$ but also a pronounced shoulder at 1855 cm^{-1} . This triplet can be assigned to the IR-active combinations of the three IR-active

Fig. 4 FT-IR spectrum of β -K₂Mo₂P₂O₁₁ in the region containing overtones. The spectrum has been amplified and corrected by baseline subtraction.

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 $(B_{1u}+B_{2u}+B_{3u})$, the four Raman-active $(A_g+B_{1g}+B_{2g}+B_{3g})$ and the inactive (A_u) Mo=O stretching modes of terminal molybdenyl groups. By using the experimentally observed maxima in the IR and Raman spectra, we can expect four maxima centred near $948+936=1884$; $948+928=1876$; $930+936=1866$; $930+928=1858$ cm⁻¹.

In the region $800-700$ cm⁻¹ we find a strong band, centred at 741 cm⁻¹, with several unresolved components in the IR spectrum and a well resolved doublet at 735 and 720 cm⁻¹ in the Raman spectrum. In this region we expect the vibrations associated to the symmetric stretch of the P–O–P diphosphate bridge $v_s(POP)$, which is known to give rise to a quite strong IR band at 740 cm^{-1} for bent-bridged diphosphates. This mode is expected to be IR-inactive for linearly-bridged diphosphate ions.²⁰ The corresponding Raman-active modes are generally of medium intensity or weak. According to the spectra, we expect one IR-active and two Raman-active components for $v_s(POP)$, (see Table 1). The assignment of the quite strong IR band at 741 cm⁻¹ to the IR-active component $v_s(POP)$ mode (B_{2u}) is in accord with the experimental P–O–P angle of 138.2°.

The assignments of the features observed at lower frequencies are far less straightforward, due to the complexity of the vibrational pattern.

In the region 720–500 cm^{-1} we can observe three strong IR bands, at 597, 557 and 513 cm^{-1} with a prominent shoulder at 613 cm^{-1} , while in the Raman pattern we cannot distinguish with certainty any peak whose intensity is definitely higher than the noise level. We can expect here the asymmetric Mo–O–Mo single bond stretches of the nearly square planar $Mo₂O₂$ units, and the asymmetric bending modes of the $PO₃$ groups. These modes should in fact give rise to weak peaks in the Raman patterns. According to Table 1 we expect 3 IR-active components for the v_{as} MoOMo modes, and six IR active modes for $\delta_{as}(\text{PO}_3)$, three of which belong to modes where the two PO₃ groups of each diphosphate unit bend "in phase" and three where they bend ''out of phase''. We can consequently tentatively assign, according to the spectra of

Table 2 Tentative assignment of IR and Raman vibrations for the β -K₂Mo₂P₂O₁₁ phase

IR $bands/cm-1$	Symmetry Mode		Symmetry	Raman $peaks/cm^{-1}$
1182 1141 1130 1113 1092	$2~\mathbf{B}_{1\mathbf{u}}$ $+2 B_{2u}$ $+2 B_{3u}$	$v_{\rm as}(\rm PO_3)$	$\rm A_g$ $\mathbf{A}_{\rm g}$	1119 1073
1066 1052 1020	$\mathbf{B}_{1\mathrm{u}}$ $+\,{\bf B}_{2{\bf u}}$ $+ B_{3u}$	$v_{\rm{sym}}(\rm{PO_3})$	$\mathbf{A}_{\rm g}$	1065
\sim 950 936 928	$B_{1u} + B_{3u}$ B_{2u} $B_{1u} + B_{3u}$	$v_{as}(P-O-P)$ $v_{sym}(Mo=O)$ v_{as} (Mo=O)	$B_{1g} + B_{3g}$ $\rm A_g$ $B_{1g} + B_{3g}$	960 948 930
741	$\mathbf{B}_{2\mathrm{u}}$	$v_{sym}(P-O-P)$	A_g+ $\mathbf{B}_{2\mathrm{g}}$	735 720
612 597	$B_{1u} + B_{2u}$ $+ B_{3u}$	$v_{\rm as}(Mo-O-Mo)$		
558 513	$B_{1u} + B_{2u}$ $+ B_{3u}$ $\mathbf{B}_{1\mathrm{u}}\!+\!\mathbf{B}_{2\mathrm{u}}$ $+ B_{3u}$	$\delta_{as}(\rm{PO}_3)$		
470	$B_{1u} + B_{2u}$ $+ B_{3u}$	$v_s/\delta(Mo-O-Mo)$ A _g		486
427 412	$B_{1u} + B_{2u}$ $+ B_{3u}$	$\delta_{\rm s}(\rm{PO}_3)$		
382 371	$B_{1u}+B_{2u}$ $+ B_{3u}$ $B_{1u} + B_{2u}$ $+ B_{3u}$	$\delta(Mo=O)$	$\mathbf{A}_{\rm g}$	375
337 323 295	B_{1u} $+\,\mathbf{B}_{2\mathrm{u}}$ $+ B_{3u}$	v_s (O-Mo-O)	$\mathbf{A}_{\rm g}$	291
282 262	$B_{1u}+B_{2u}$ $+ B_{3u}$ $\mathbf{B}_{1\mathrm{u}}\!+\!\mathbf{B}_{2\mathrm{u}}$ $+ B_{3u}$	rock $PO3$	$\rm A_g$	267
245		twist PO ₃	$\rm A_g$	247
237 217	B_{2u}	rock Mo-O-Mo δ (POP)	$\mathbf{A}_{\rm g}$	207
194 173 153 127		KO ₉ modes and lattice modes	A_g	171
111 102 81 69				

other diphosphates and of similar compounds, the complex band near 600 cm⁻¹ to v_{as} (MoOMo) modes and the bands at 558 and 513 cm⁻¹ to $\delta_{as}(\text{PO}_3)$ modes.

In the region 500–450 cm^{-1} we found a quite strong IR band at 471 cm^{-1} and a weak Raman peak at 486 cm^{-1}. These features can be assigned to the symmetric stretches of the Mo– O–Mo bridges (this mode is expected to give rise to quite intense features in Raman as well as in IR).

The IR spectrum shows a complex absorption in the range $450-350$ cm⁻¹. In this region a medium-strong Raman peak is found at 375 cm^{-1} , that can correspond to a sharp IR component at 382 cm⁻¹, which also shows a shoulder at 371 cm⁻¹. Possible

assignments are to the bending modes of terminal $Mo=O$ bonds, that is expected to give in IR three modes when the two molybdenyls of every unit Mo2O4 bend in their average plane and three when they bend out-of-plane. The two IR components at 427 and 412 cm^{-1} , that do not have detectable counterparts in the Raman pattern, can be assigned to the symmetric deformations of the diphosphate groups $\delta_s(\text{PO}_3)$.

In the range $350-200 \text{ cm}^{-1}$ another very complex absorption is found in the IR spectrum, where also several Raman peaks fall. The main peak maxima are in IR at 295 and 262 cm^{-1} and in Raman at 291 and 267 cm^{-1} . These components can be assigned, in both techniques, to different deformations of the PO_3 groups of diphosphates, *i.e.* the twisting and the rocking modes. A further sharp peak in Raman is found at 207 cm⁻ that can correspond to a quite sharp IR band at 194 cm^{-1} and can be tentatively assigned to the deformation of the P–O–P diphosphate bridges.

The modes arising from the vibrations of the $KO₉$ polyhedra are expected to lie below 190 cm^{-1} , together with modes that have the predominant character of lattice modes (translational and rotational modes of the P_2O_7 and M_2O_4 molecular units).

2. The vibrational spectra of α -K₂Mo₂P₂O₁₁

The IR spectrum of α -K₂Mo₂P₂O₁₁ is shown in Fig. 1b. We could not obtain a high-quality FT-Raman spectrum for this phase, due to thermoluminescence effects. The Raman spectrum we recorded (not shown here because of the high noise level) only shows definite peaks at 940 and 813 cm^{-1} . The tentative assignment of the vibrational modes is given in Table 3.

This phase crystallises into the space group $P2_1/c = C_{2h}^5 =$ no. 14, with $Z=4$. As mentioned above, it is a monophosphate and its formula can thus be written as $K_2Mo_2O_3(PO_4)_2$.^{11,15} The molybdenum atom is in octahedral coordination characteristic of $Mo(v), i.e.$ with one free apex corresponding to one short $Mo=O$ molybdenum bond (1.683 Å) . Four of the five other apices are linked to four PO_4 tetrahedra and the fifth is shared with a similar $MoO₆ octahedron, forming Mo₂O₁₁ bioctahedral units in which$ the two molybdenyl bonds are directed along the same direction. The Mo–O–Mo angle in the $Mo₂O₁₁$ unit is 149.79°.^{11,15}

As already discussed in the case of the β phase, the complex IR region above 800 cm^{-1} (Fig. 1, curve b) contains symmetric and asymmetric $P=O$ and $Mo=O$ stretches. In particular, the 950–1100 cm⁻¹ spectral range shows PO_x terminal stretches of phosphate groups. In the α phase the average P–O bond length is about 1.55 Å, while in the β structure it ranges between 1.50 and 1.52 Å, so being definitively shorter. This causes a decrease in the wavenumbers of the v_{as} (1094, 1073, \sim 1035, \sim 1011 and 991 cm⁻¹) and $v_s(PO_3)$ (975 and 965 cm⁻¹) modes. As we can argue from Fig. 3b, the $v(POP)$ vibration cannot be taken into account in this compound, since the framework contains only monophosphate units. The Raman band at 940 cm^{-1} can be assigned to the terminal molybdenyl groups $Mo=O$ symmetric stretching mode, while two of the 4 IR-active corresponding asymmetric modes give rise to the bands at 927 and 915 cm^{-1} (spectrum 1b);the shift towards lower frequencies of suchpeaks in comparison with the above discussed structure, can be ascribed to the Mo=O bond length, which is slightly shorter in the latter case.

The two bands at 819 in IR and 813 cm^{-1} in Raman (absent in the β structure) could likely be due to IR- and Raman-active components of the asymmetric stretches Mo–O–Mo $(\text{angle}=149.79^{\circ})$, which are definitely shorter here with respect to the β phase.²⁵

Between 700 and 500 cm⁻¹ we can expect the asymmetric Mo–O–Mo single bond stretches (650 and 615 cm⁻¹) and the asymmetric bending modes of the $PO₃$ groups (576, 556 and 538 cm^{-1}). The corresponding symmetric stretch of the Mo–O– Mo bridges should be represented by the peak at 471 cm⁻¹. A possible assignment of the peak at 440 cm^{-1} is to the symmetric $\delta_s(PO_3)$ bending of the phosphate groups, while the bands

\mathbb{R} 201021 2011 phase			Culvior O_6 phase				
IR bands/ cm^{-1}	Tentative assignment	Raman peaks/ cm^{-1}	IR bands/ cm^{-1}	Tentative assignment	Raman peaks/ cm^{-1}		
1094 1073 1035 1011 991	$v_{\rm as}(\rm PO_3)$		\sim 1093 \sim 1047 1029 $\sim\!1021$ 1005	$v_{\rm as}(\rm PO_3)$	1106 1074		
975 965	$v_s(PO_3)$		977 946	$v_s(PO_3)$ v_{as} (Mo=O)	988		
	$v_s(Mo=O)$	940	936	$v_s(Mo=O)$	927		
927 915	$v_{\rm as}$ (Mo=O)		\sim 920-790 761 \sim 750	$v(P-O-Mo)$	803		
819 650	$v_{\rm as}$ Mo-O-Mo	813	$\sim\!725$ 631		722		
615			~ 622	$v_{as}(Mo-O-Mo)$			
576 556 538	$\delta_{as}(\rm{PO}_3)$		577 541 530	δ_{as})PO ₃)			
514 490			482 \sim 477	$v_s(Mo-O-Mo)$	\sim 480		
471 440	$v_s(Mo-O-Mo)$ $\delta_{\rm s}(\rm{PO}_3)$		435 416	δ _s (PO ₃)	412		
414 404	$\delta(Mo=O)$		373 359 350	$\delta(Mo=O)$	380		
364 355 328 287			339 330	v_s (O-Mo-O)	\sim 359		
284 278 269 250			\sim 301 \sim 287 $\sim\!280$ 272	rock PO ₃	267		
244 221 217 203			252 245	twist $PO3$	228		
198 191 179 173			224 212 206 $\sim\!201$	rock Mo-O-Mo			
168 154 134 119			169 \sim 164 144 129 100 77	$CdO6$ and lattice modes	155		

Table 3 Tentative assignment of IR and Raman vibrations for the α - K_2M_0 ¹⁰¹

Table 4 Tentative assignment of IR and Raman vibrations for the $CdMoPO₆ phase$

around 400 cm^{-1} are assignable to the bending modes of terminal Mo=O bonds. Other deformations of $PO₃$ units are responsible for the bands between 350 and 200 cm^{-1} , while the features below 200 cm⁻¹ could be taken into account to explain vibrations of KO₉ polyhedra and lattice modes.

3. The vibrational spectra of $CdMoPO₆$

The monophosphate containing $Mo(v)$ CdMoPO₆^{11,17} crystallises in the $P2_1/c$ space group (C_{2h}^5 , no. 14), with $Z=4$. As in the two previous structures, the molybdenum atom presents the characteristic octahedral coordination of Mo(V) with one short Mo=O bond (1.693 Å), corresponding to a molybdenyl group. Each octahedron thus presents one free corner and shares three apices with three $PO₄$ tetrahedra and one edge with another $MoO₆ octahedron, leading to Mo₂O₁₀ bioctahedral units (Mo–$ O–Mo angle = 84.05° , similar to that encountered in the β - $K_2Mo_2P_2O_{11}$ compound). These Mo_2O_{10} units are similar to those described for β -K₂Mo₂P₂O₁₁, except that in $CdMoO₂(PO₄)$ the two molybdenyl bonds of the units are antiparallel (whereas they are parallel in β -K₂Mo₂P₂O₁₁), *i.e.* they are in a *trans* symmetry.

As shown in Fig. 3c, $CdMoPO₆$ can be regarded as constituted by an intermediate surrounding between that of α - and β -K₂Mo₂P₂O₁₁, *i.e.* two units of two edge-sharing molybdenum octahedra (like for β -K₂Mo₂P₂O₁₁) linked by monophosphate groups (as in the case of α -K₂Mo₂P₂O₁₁). For this reason we try here to assign IR and Raman modes by correlation with the two previous structures.

The FT-IR and Raman spectra of $CdMoPO₆$ are respectively shown in Figs. 1c and 2c. The tentative assignment of the vibrational modes is given in Table 4. As already discussed for the parent compounds, the complex feature between 1100 and 1000 cm^{-1} in the IR spectrum should be assigned to the asymmetric stretches of $PO₃$ monophosphate groups, with the corresponding Raman peaks situated at 1106 and 1074 cm^{-1} . The symmetric modes arising from the same units could be

found at 977 (IR) and 988 (Raman) cm^{-1} . The rather intense peak at 946 cm^{-1} in the IR spectrum can be likely attributed to the asymmetric stretch of the molybdenyl bonds, the corresponding symmetric vibration giving rise to the main features at 936 (IR) and 927 (Raman) cm^{-1} .

The broad bump characterising the IR spectrum between \sim 920 and 790 cm⁻¹ is rather difficult to interpret. Being at too low wavenumbers (and too intense) for an electronic effect, it could be assigned to the P–O–Mo stretches, together with the 761, \sim 750 and \sim 725 cm⁻¹ IR bands and the 803 and 722 cm^{-1} Raman peaks. These features are in fact in a similar vibrational region to P–O–P bands for diphosphates, but at higher wavenumbers (Raman) and less resolved (IR) due to the average longer Mo–O bonds (from 1.944 to 2.130 Å) with respect to P–O bonds $(1.518-1.559 \text{ Å})$. The IR band at 631 cm⁻¹ and the shoulder at 622 cm⁻¹ can be assigned to the asymmetric stretching mode v_{as} (Mo–O–Mo), while the corresponding symmetric features are suggested at 482 and \sim 477 cm⁻¹ (IR) and \sim 480 cm⁻¹ (Raman). The triplet at 577, 541 and 530 cm⁻¹ in the IR spectrum can be due to the $\delta_{as}(\overline{P}O_3)$. The symmetric bending $\delta_{s}(\overline{P}O_3)$ can be attributed to the vibrations at 435 and 416 cm⁻¹ (IR) and at 412 cm⁻¹ (Raman). The deformation $\delta(Mo=O)$ of the molybdenyl bond could give rise to the bands at 373, 359 and 350 cm⁻¹ in the IR spectrum and to the Raman peak at 380 cm^{-1} . The doublet at 339 and 330 cm⁻¹ in IR and the band at 359 cm⁻¹ in Raman have been assigned to the v_s (O–Mo–O). The features between 300 and 260 cm^{-1} can be due to the rocking mode of PO₃ units, whereas the corresponding twisting should give rise to bands at 252 and 245 cm⁻¹ in IR and 228 cm⁻¹ in Raman. Between 225 and 200 cm^{-1} in the IR spectrum may occur the rocking mode for the Mo–O–Mo bonds. Features below this are assignable to $CdO₆$ and lattice modes.

Conclusion

The analysis of the vibrational modes reported above with details for α and β phases and the corresponding interpretation of the experimental IR and Raman spectra allow us to confirm the usefulness of vibrational spectroscopy for structural characterisation of oxides and salts.

This is well known and indeed has been already discussed many times for simple metal phosphates. In the present case the spectral analysis of more complex mixed K–Mo phosphates allow us to obtain criteria for the spectra–structure relation of complex mixed Mo–O species.

In the case of the β -phase, the pairs of edge-sharing octahedra present nearly parallel short molybdenyl bonds in a cis configuration. They are characterised by three IR active modes and by three Raman active modes, giving rise to doublets in both spectra. The splitting of the absorption is between 10 and 20 cm^{-1} and the symmetric stretching modes are detected at higher frequencies than the asymmetric ones. By contrast the α polymorph presents bioctahedral units shared by corners; again two short molybdenyl bonds are present, being nearly parallel in the cis configuration. In this case we can observe a strong Raman peak at 940 cm^{-1} , which is confidently assigned to the asymmetric stretching mode, while the split IR mode is at lower frequencies. Again the splitting between symmetric and asymmetric stretching modes is in the range between 10 and 25 cm^{-1} . This confirms that molybdenyl(v) species in *cis* configuration, either edge- or corner-sharing, give rise to a couple of bonds in the region 950–910 cm^{-1} , the higher frequency being assigned to the symmetric stretching mode, and stronger in the Raman spectrum.

For $CdMoPO₆$, edge-sharing pairs of Mo octahedra exist

with molybdenyl bonds in *trans*-configuration. Also in this case the splitting between symmetric and asymmetric modes is 10 cm^{-1} or higher, but now the symmetric stretch (the most intense in Raman spectroscopy) is at lower frequencies. Raman spectroscopy has been largely used to investigate the structures of amorphous materials such as glasses and of otherwise undetectable species such as supported species on catalysts. In particular^{14,26–29} Raman and infrared spectroscopies are useful to characterise monomeric molybdenum species (where only one IR and Raman bands are expected) from dimeric species (where at least two bands are expected in the spectra). This suggests that vibrational spectroscopies and, in particular, the Raman technique are suitable to distinguish different molybdenum coordinate structures inserted in a matrix or coated on a support.

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